

Spatial decay of rotating waves in parabolic systems

Nonlinear Waves, CRC 701, Bielefeld, June 19, 2013

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CRC 701

Outline

- 1 Introduction: Rotating pattern in \mathbb{R}^d
- 2 Main result: Exponential decay of v_\star
- 3 Outline of proof: Exponential decay of v_\star
- 4 Spectrum of rotating waves

Rotating Patterns in \mathbb{R}^d

Consider a **reaction diffusion system**

$$(1) \quad \begin{aligned} u_t(x, t) &= A\Delta u(x, t) + f(u(x, t)), & t > 0, x \in \mathbb{R}^d, d \geq 2, \\ u(x, 0) &= u_0(x) & , t = 0, x \in \mathbb{R}^d. \end{aligned}$$

where $u : \mathbb{R}^d \times [0, \infty[\rightarrow \mathbb{R}^N$, $A \in \mathbb{R}^{N,N}$, $f \in C^2(\mathbb{R}^N, \mathbb{R}^N)$.

Assume a **rotating wave** solution $u_* : \mathbb{R}^d \times [0, \infty[\rightarrow \mathbb{R}^N$ of (1)

$$u_*(x, t) = v_*(e^{-tS}x)$$

$v_* : \mathbb{R}^d \rightarrow \mathbb{R}^N$ profile (pattern), $0 \neq S \in \mathbb{R}^{d,d}$ skew-symmetric.

Transformation (into a rotating frame): $v(x, t) = u(e^{tS}x, t)$ solves

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$$\langle Sx, \nabla v(x) \rangle := \sum_{i=1}^d \sum_{j=1}^d S_{ij}x_j D_i v(x) \quad (\text{drift term}).$$

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$$\langle Sx, \nabla v(x) \rangle := \sum_{i=1}^{d-1} \sum_{j=i+1}^d S_{ij} (x_j D_i - x_i D_j) v(x) \quad (\text{rotational term}).$$

Rotating Patterns in \mathbb{R}^d

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v_* is a stationary solution of (2).

Question: How to show exponential decay of v_* at $|x| = \infty$?

Consequence: Exponentially small error by restriction to bounded domain.

Rotating Patterns in \mathbb{R}^d

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$d = 2$: Spectral stability implies nonlinear stability.

 [BL] W.-J. Beyn, J. Lorenz. 2008.

Example

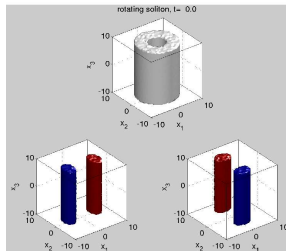
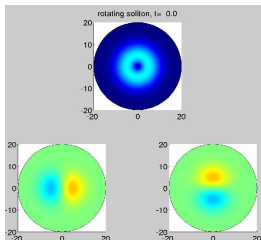
Consider the **quintic complex Ginzburg-Landau equation** (QCGL):

$$u_t = \alpha \Delta u + u \left(\mu + \beta |u|^2 + \gamma |u|^4 \right), \quad u = u(x, t) \in \mathbb{C}$$

with $u : \mathbb{R}^d \times [0, \infty[\rightarrow \mathbb{C}$, $d \in \{2, 3\}$. For the parameters

$$\alpha = \frac{1}{2} + \frac{1}{2}i, \beta = \frac{5}{2} + i, \gamma = -1 - \frac{1}{10}i, \mu = -\frac{1}{2}$$

this equation exhibits so called **spinning soliton** solutions.



Applications: superconductivity, superfluidity, nonlinear optical systems.

Example

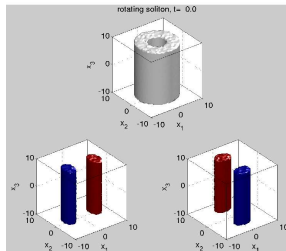
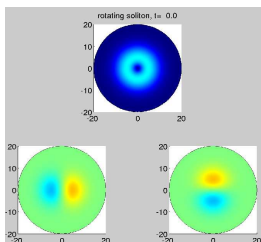
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[CMM] L.-C. Crasovan, B.A. Malomed, D. Mihalache. 2001

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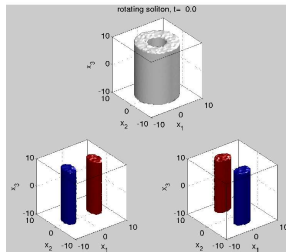
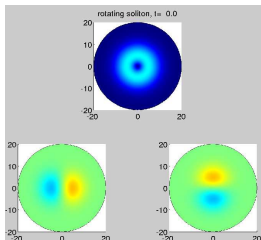
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Main result: Exponential decay of v_\star

Theorem: (Exponential Decay of v_\star)

Let $f(v_\infty) = 0$ and $\operatorname{Re} \sigma(Df(v_\infty)) < 0$. Under further assumptions holds:
For every $1 < p < \infty$, $0 < \vartheta < 1$ and for every radially nondecreasing weight function $\theta \in C(\mathbb{R}^d, \mathbb{R})$ of exponential growth rate $\eta \geq 0$ with

$$0 \leq \eta^2 \leq \vartheta \frac{2}{3} \frac{a_0 b_0}{a_{\max}^2 p^2}$$

there exists $K_1 = K_1(A, f, v_\infty, d, p, \theta, \vartheta) > 0$ with the following property:
Every classical solution v_\star of

$$A \Delta v(x) + \langle Sx, \nabla v(x) \rangle + f(v(x)) = 0, \quad x \in \mathbb{R}^d,$$

such that $v_\star - v_\infty \in L^p(\mathbb{R}^d, \mathbb{R}^N)$ and

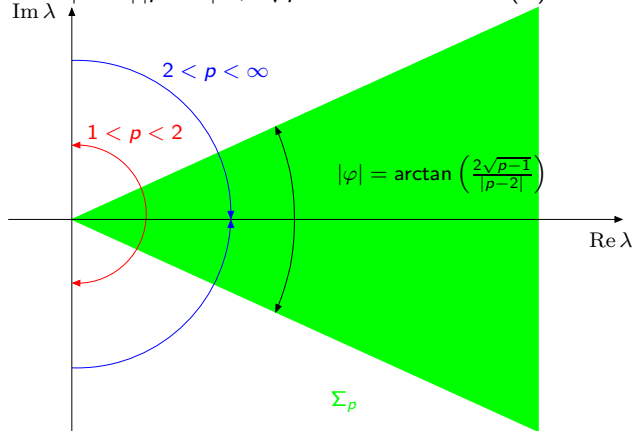
$$\sup_{|x| \geq R_0} |v_\star(x) - v_\infty| \leq K_1 \text{ for some } R_0 > 0$$

satisfies

$$v_\star - v_\infty \in W_\theta^{1,p}(\mathbb{R}^d, \mathbb{R}^N) \text{ (weighted Sobolev space).}$$

Main result: The assumptions

- $\operatorname{Re} \lambda > 0$ and $\operatorname{Im} \lambda$ $|\operatorname{Im} \lambda| |\rho - 2| \leq 2\sqrt{\rho - 1} \operatorname{Re} \lambda \quad \forall \lambda \in \sigma(A)$



$A, Df(v_\infty) \in \mathbb{R}^{N,N}$ simultaneously diagonalizable over \mathbb{C}

- $a_0 \leq \operatorname{Re} \lambda, \quad |\lambda| \leq a_{\max} \quad \forall \lambda \in \sigma(A)$
 $\operatorname{Re} \mu \leq -b_0 < 0 \quad \forall \mu \in \sigma(Df(v_\infty))$

Main result: The assumptions

- A positive function $\theta \in C(\mathbb{R}^d, \mathbb{R})$ is called a **weight function of exponential growth rate** $\eta \geq 0$ provided that

$$\exists C_\theta > 0 : \theta(x+y) \leq C_\theta \theta(x) e^{\eta|y|} \quad \forall x, y \in \mathbb{R}^d.$$

 [ZM] S. Zelik, A. Mielke. 2009.

Examples: $\mu \in \mathbb{R}$, $x \in \mathbb{R}^d$

$$\theta_1(x) = \exp(-\mu|x|), \quad \theta_3(x) = \exp\left(-\mu\sqrt{|x|^2+1}\right),$$

$$\theta_2(x) = \cosh(\mu|x|), \quad \theta_4(x) = \cosh\left(\mu\sqrt{|x|^2+1}\right).$$

- **Exponentially weighted Sobolev spaces:** $1 \leq p \leq \infty$, $k \in \mathbb{N}_0$

$$L_\theta^p(\mathbb{R}^d, \mathbb{R}^N) := \{v \in L_{loc}^1(\mathbb{R}^d, \mathbb{R}^N) \mid \|\theta v\|_{L^p} < \infty\},$$

$$W_\theta^{k,p}(\mathbb{R}^d, \mathbb{R}^N) := \{v \in L_\theta^p(\mathbb{R}^d, \mathbb{R}^N) \mid D^\beta u \in L_\theta^p(\mathbb{R}^d, \mathbb{R}^N) \forall |\beta| \leq k\}.$$

Outline of proof: Exponential Decay of v_*

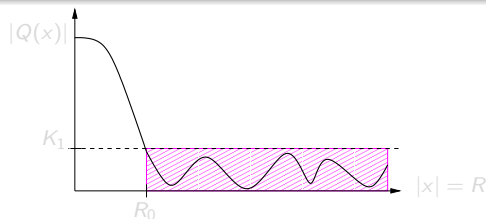
Consider the nonlinear problem

$$A\Delta v_*(x) + \langle Sx, \nabla v_*(x) \rangle + f(v_*(x)) = 0, \quad x \in \mathbb{R}^d, \quad d \geq 2.$$

1. **Far-Field Linearization:** $f \in C^1$, **Taylor's theorem**, $f(v_\infty) = 0$

$$a(x) = \int_0^1 Df(v_\infty + t(v_*(x) - v_\infty)) dt, \quad w(x) := v_*(x) - v_\infty$$

$$A\Delta w(x) + \langle Sx, \nabla w(x) \rangle + a(x)w(x) = 0, \quad x \in \mathbb{R}^d.$$



Outline of proof: Exponential Decay of v_\star

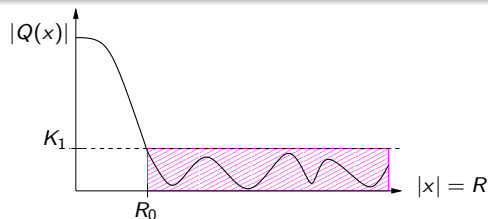
Consider the nonlinear problem

$$A\Delta v_\star(x) + \langle Sx, \nabla v_\star(x) \rangle + f(v_\star(x)) = 0, \quad x \in \mathbb{R}^d, \quad d \geq 2.$$

2. Decomposition of a :

$$Df(v_\infty) + Q(x) = \int_0^1 Df(v_\infty + t(v_\star(x) - v_\infty)) dt, \quad w(x) := v_\star(x) - v_\infty$$

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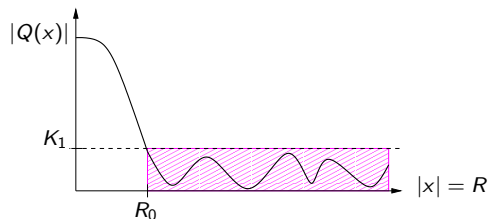
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3. Decomposition of Q :

$$\begin{aligned} Q(x) &= Q_\varepsilon(x) + Q_c(x), \\ Q, Q_\varepsilon, Q_c &\in L^\infty(\mathbb{R}^d, \mathbb{R}^{N,N}), \\ Q_\varepsilon \text{ small, i.e. } &\|Q_\varepsilon\|_{L^\infty} < K_1, \\ Q_c &\text{ compactly supported.} \end{aligned}$$

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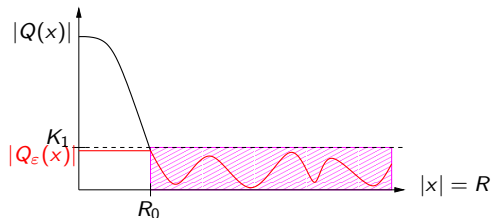
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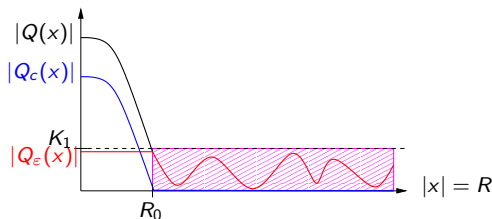
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Exponential Decay: To show **exponential decay** for the solution v_* of

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investigate the **far-field linearization** (w.l.o.g. $v_\infty = 0$)

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Operators: Study the following operators

$$\mathcal{L}_Q v := A\Delta v + \langle S \cdot, \nabla v \rangle + Df(v_\infty)v + Q_\varepsilon v + Q_c v,$$

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$$\mathcal{L}_\infty v := A\Delta v + \langle S \cdot, \nabla v \rangle + Df(v_\infty)v,$$

$$\mathcal{L}_0 v := A\Delta v + \langle S \cdot, \nabla v \rangle \quad (\text{Ornstein-Uhlenbeck operator}). \quad (\text{max. domain})$$

Ornstein-Uhlenbeck Operator

Let $P, B \in \mathbb{R}^{d,d}$, $P = P^T$, $P > 0$ and $B \neq 0$.

$$\nabla^T P \nabla v(x) + \langle Bx, \nabla v(x) \rangle = \sum_{i=1}^d \sum_{j=1}^d D_i (P_{ij} D_j v(x)) + \sum_{i=1}^d \sum_{j=1}^d D_i v(x) B_{ij} x_j, \quad x \in \mathbb{R}^d$$

Here: $P = I_d$ and $B = S$.

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Operators: Study the following operators

$$\mathcal{L}_Q v := A\Delta v + \langle S \cdot, \nabla v \rangle + Df(v_\infty)v + Q_\varepsilon v + Q_c v,$$

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$$\mathcal{L}_\infty v := A\Delta v + \langle S \cdot, \nabla v \rangle + Df(v_\infty)v, \quad (\text{exp. decay})$$

$$\mathcal{L}_0 v := A\Delta v + \langle S \cdot, \nabla v \rangle \quad (\text{Ornstein-Uhlenbeck operator}). \quad (\text{max. domain})$$

Ornstein-Uhlenbeck Operator

Let $P, B \in \mathbb{R}^{d,d}$, $P = P^T$, $P > 0$ and $B \neq 0$.

$$\nabla^T P \nabla v(x) + \langle Bx, \nabla v(x) \rangle = \sum_{i=1}^d \sum_{j=1}^d D_i (P_{ij} D_j v(x)) + \sum_{i=1}^d \sum_{j=1}^d D_i v(x) B_{ij} x_j, \quad x \in \mathbb{R}^d$$

Here: $P = I_d$ and $B = S$.

Exponential Decay: To show **exponential decay** for the solution v_* of

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 [MPV] G. Metafune, D. Pallara, V. Vespri. 2005.

 [MPRS] G. Metafune. 2001.

The operator \mathcal{L}_0

Ornstein-Uhlenbeck operator

$$[\mathcal{L}_0 v](x) = A\Delta v(x) + \langle Sx, \nabla v(x) \rangle, \quad x \in \mathbb{R}^d, \quad d \geq 2.$$

↓

$$H_0(x, \xi, t) = (4\pi tA)^{-\frac{d}{2}} \exp\left(- (4tA)^{-1} \left| e^{tS} x - \xi \right|^2\right), \quad x, \xi \in \mathbb{R}^d, \quad t > 0.$$

↓

Semigroup in $L^p(\mathbb{R}^d, \mathbb{R}^N)$, $1 \leq p \leq \infty$

$$[T_0(t)v](x) = \int_{\mathbb{R}^d} H_0(x, \xi, t)v(\xi)d\xi, \quad t > 0.$$

strong ↓ continuity

Infinitesimal generator

$$(A_p, \mathcal{D}(A_p)), \quad 1 \leq p < \infty.$$

semigroup theory ✓

A-priori ↓ estimates

↘ \mathcal{L}_0 : L^p -resolvent est.

unique solv. of
resolvent equ.,

$$1 \leq p < \infty$$

exponential
decay,

$$1 \leq p < \infty$$

max. domain and
max. realization,

$$1 < p < \infty$$

$$(\lambda I - A_p)v_* = g \in L^p.$$

$$v_* \in W_{\theta}^{1,p}.$$

$$A_p = \mathcal{L}_0 \text{ on } \mathcal{D}(A_p) = \mathcal{D}^p(\mathcal{L}_0).$$

Spectrum of localized rotating waves

Linearization about the profile of the rotating wave

$$[\mathcal{L}v](x) = A\Delta v(x) + \langle Sx, \nabla v(x) \rangle + Df(v_*(x))v(x), \quad x \in \mathbb{R}^d, \quad d \geq 2.$$

Eigenvalue problem

$$[\mathcal{L}v](x) = \lambda v(x), \quad x \in \mathbb{R}^d, \quad d \geq 2, \quad \lambda \in \mathbb{C}.$$

A **rotating wave** solution $u_*(x, t) = v_*(e^{-tS}x)$ is **spectrally stable** if

$$\sigma(\mathcal{L}) \subseteq \{\lambda \in \mathbb{C} \mid \operatorname{Re} \lambda \leq 0\}.$$

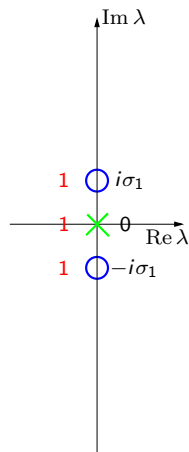
Decompose the **spectrum** $\sigma(\mathcal{L})$ into

$$\sigma(\mathcal{L}) = \sigma_{\text{ess}}(\mathcal{L}) \dot{\cup} \sigma_{\text{pt}}(\mathcal{L}),$$

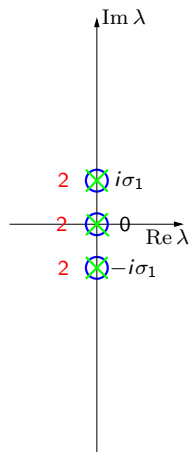
with essential spectrum $\sigma_{\text{ess}}(\mathcal{L})$ and point spectrum $\sigma_{\text{pt}}(\mathcal{L})$.

Illustration: Point spectrum of \mathcal{L}

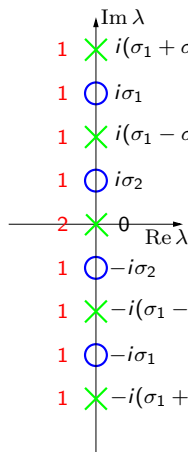
$\lambda \in (\sigma(S) \cup \{\lambda + \mu \mid \lambda, \mu \in \sigma(S), \lambda \neq \mu\}) \subseteq \sigma_{\text{pt}}(\mathcal{L})$ with algebraic multiplicity



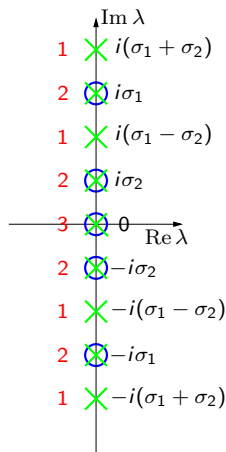
$d = 2$
 $\dim \text{SE}(2) = 3$



$d = 3$
 $\dim \text{SE}(3) = 6$



$d = 4$
 $\dim \text{SE}(4) = 10$



$d = 5$
 $\dim \text{SE}(5) = 15$

Point spectrum of \mathcal{L}

Theorem: (Point spectrum of \mathcal{L})

Let $A \in \mathbb{R}^{N,N}$ satisfy the main assumptions, $f \in C^2(\mathbb{R}^N, \mathbb{R}^N)$, $1 < p < \infty$ and let v_* be a classical solution of $[\mathcal{L}_0 v](x) + f(v(x)) = 0$. Then

$$v(x) = \langle C^{rot}x + I_d C^{tra}, \nabla v_*(x) \rangle, \quad x \in \mathbb{R}^d, \quad C^{rot} \in \mathfrak{so}(d), \quad C^{tra} \in \mathbb{R}^d$$

solves $\mathcal{L}v = \lambda v$, whenever $(\lambda, (C^{rot}, C^{tra}))$ solves

$$\begin{aligned}\lambda C^{rot} &= -SC^{rot} + (SC^{rot})^T, \\ \lambda C^{tra} &= -SC^{tra}.\end{aligned}$$

In particular:

- $\sigma(S) \cup \{\lambda + \mu \mid \lambda, \mu \in \sigma(S), \lambda \neq \mu\} \subseteq \sigma_{\text{pt}}(\mathcal{L})$,
- $v(x) = \langle Sx, \nabla v_*(x) \rangle$ eigenfunction of $\lambda = 0$ for every $d \geq 2$,
- **group action** $\Rightarrow \sigma_{\text{pt}}(\mathcal{L})$,
- Theorem also valid for spiral waves, scroll waves, scroll rings.

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- Theorem also valid for spiral waves, scroll waves, scroll rings.

Exponential decay of v

Theorem: (Exponential decay of v)

Let the assumptions of the main result be satisfied.

The every classical solution $v \in L^p(\mathbb{R}^d, \mathbb{C}^N)$ of

$$A\Delta v(x) + \langle Sx, \nabla v(x) \rangle + Df(v_*(x))v(x) = \lambda v(x), \quad x \in \mathbb{R}^d,$$

with $\lambda \in \mathbb{C}$ and $\operatorname{Re} \lambda \geq -\frac{b_0}{3}$ satisfies

$$v \in W_{\theta}^{1,p}(\mathbb{R}^d, \mathbb{C}^N)$$

- v_* exp. localized $\Rightarrow v$ exp. localized (with same rate)

Outline of proof: Point spectrum of \mathcal{L}

$$(3) \quad 0 = A\Delta v_*(x) + \langle Sx, \nabla v_*(x) \rangle + f(v_*(x)), \quad x \in \mathbb{R}^d, \quad d \geq 2.$$

1. Group action: Apply $a(R, \tau)$ to (3)

$$0 = a(R, \tau) (A\Delta v_*(x) + \langle Sx, \nabla v_*(x) \rangle + f(v_*(x)))$$

2. Derivative $\frac{d}{d(R, \tau)}$ at $(R, \tau) = (I_d, 0)$ leads to $\frac{d(d+1)}{2}$ equations

$$0 = (x_j D_i - x_i D_j) (A\Delta v_*(x) + \langle Sx, \nabla v_*(x) \rangle + f(v_*(x)))$$

$$0 = D_l (A\Delta v_*(x) + \langle Sx, \nabla v_*(x) \rangle + f(v_*(x)))$$

for $i = 1, \dots, d-1, j = i+1, \dots, d, l = 1, \dots, d$.

3. Commutator relations for differential operators yield, $D^{(ij)} := x_j D_i - x_i D_j$

$$0 = \mathcal{L} \left(D^{(ij)} v_*(x) \right) + \sum_{\substack{n=1 \\ n \neq j}}^d S_{in} D^{(jn)} v_*(x) - \sum_{\substack{n=1 \\ n \neq i}}^d S_{jn} D^{(in)} v_*(x)$$

$$0 = \mathcal{L} (D_l v_*(x)) - \sum_{n=1}^d S_{ln} D_n v_*(x)$$

Outline of proof: Point spectrum of \mathcal{L}

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4. **Finite-dimensional eigenvalue problem:** The ansatz

$$v(x) = \sum_{i=1}^{d-1} \sum_{j=i+1}^d C_{ij}^{rot} (x_j D_i - x_i D_j) v_\star(x) + \sum_{l=1}^d C_l^{tra} D_l v_\star(x), \quad C_{ij}^{rot}, C_l^{tra} \in \mathbb{C}$$

reduces $\mathcal{L}v = \lambda v$ to a

$$\begin{aligned} \lambda C^{rot} &= -S C^{rot} + (S C^{rot})^T, \\ \lambda C^{tra} &= -S C^{tra}. \end{aligned}$$

Note: S is unitary diagonalizable.

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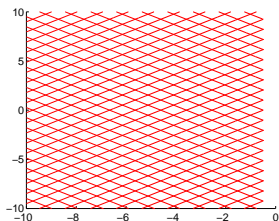
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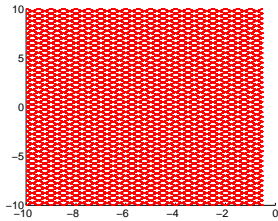
Note: S is unitary diagonalizable.

Illustration: Essential spectrum of \mathcal{L}

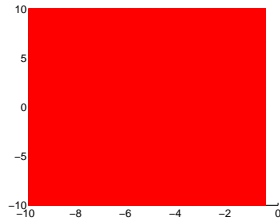
$$\left\{ -\lambda(\omega) + i \sum_{l=1}^k n_l \sigma_l \mid \lambda(\omega) \text{ eigenvalue of } \omega^2 A - Df(v_\infty) \right\} \subseteq \sigma_{\text{ess}}(\mathcal{L})$$



$d = 2$ or 3



$d = 4$ (not dense)



$d = 4$ (dense)

Essential spectrum of \mathcal{L}

Theorem: (Essential spectrum of ν)

Let the assumptions of the main result be satisfied. Moreover, let $\pm i\sigma_1, \dots, \pm i\sigma_k$ denote the nonzero eigenvalues of S and let $\lambda(\omega)$ denote an eigenvalue of $\omega^2 A - Df(v_\infty)$ for some $\omega \in \mathbb{R}$. Then

$$\left\{ \lambda = -\lambda(\omega) - i \sum_{l=1}^k n_l \sigma_l \in \mathbb{C} \mid n_l \in \mathbb{Z}, \omega \in \mathbb{R} \right\} \subseteq \sigma_{\text{ess}}(\mathcal{L})$$

in $L^p(\mathbb{R}^d, \mathbb{C}^N)$.

- Far-field linearization $\Rightarrow \sigma_{\text{ess}}(\mathcal{L})$

Dispersion relation: $\lambda \in \sigma_{\text{ess}}(\mathcal{L})$ if

$$\det \left(\lambda I_N + \omega^2 A + i \sum_{l=1}^k n_l \sigma_l I_N - Df(v_\infty) \right) = 0 \text{ for some } \omega \in \mathbb{R}.$$

Outline of proof: Essential spectrum of \mathcal{L}

Linearization at the profile v_* :

$$[\mathcal{L}v](x) = A\Delta v(x) + \langle Sx, \nabla v(x) \rangle + Df(v_\infty)v(x) + Q(x)v(x)$$

$$Q(x) := Df(v_*(x)) - Df(v_\infty), \quad \sup_{|x| \geq R} |Q(x)|_2 \rightarrow 0 \text{ as } R \rightarrow \infty$$

1. Orthogonal transformation: $S \in \mathbb{R}^{d,d}$, $S^T = -S$, $S = P\Lambda_{\text{block}}^S P^T$.

$T_1(x) = Px$ yields

$$[\mathcal{L}_1v](x) = A\Delta v(x) + \langle \Lambda_{\text{block}}^S x, \nabla v(x) \rangle + Df(v_\infty)v(x) + Q(T_1(x))v(x)$$

with

$$\langle \Lambda_{\text{block}}^S x, \nabla v(x) \rangle = \sum_{l=1}^k \sigma_l (x_{2l} D_{2l-1} - x_{2l-1} D_{2l}) v(x).$$

Outline of proof: Essential spectrum of \mathcal{L}

Orthogonal transformation:

$$[\mathcal{L}_1 v](x) = A\Delta v(x) + \langle \Lambda_{\text{block}}^S x, \nabla v(x) \rangle + Df(v_\infty)v(x) + Q(T_1(x))v(x)$$

$$\langle \Lambda_{\text{block}}^S x, \nabla v(x) \rangle = \sum_{l=1}^k \sigma_l (x_{2l} D_{2l-1} - x_{2l-1} D_{2l}) v(x)$$

2. Several planar polar coordinates: Transformation

$$\begin{pmatrix} x_{2l-1} \\ x_{2l} \end{pmatrix} = T(r_l, \phi_l) := \begin{pmatrix} r_l \cos \phi_l \\ r_l \sin \phi_l \end{pmatrix}, \quad l = 1, \dots, k, \quad \phi_l \in]-\pi, \pi], \quad r_l > 0.$$

yields for $\xi = (r_1, \phi_1, \dots, r_k, \phi_k, x_{2k+1}, \dots, x_d)$ with total transformation $T_2(\xi)$,
 $Q(\xi) := Q(T_1(T_2(\xi)))$

$$[\mathcal{L}_2 v](x) = A \left[\sum_{l=1}^k \left(\partial_{r_l}^2 + \frac{1}{r_l} \partial_{r_l} + \frac{1}{r_l^2} \partial_{\phi_l}^2 \right) + \sum_{l=2k+1}^d \partial_{x_l}^2 \right] v(\xi) \\ - \sum_{l=1}^k \sigma_l \partial_{\phi_l} v(\xi) + Df(v_\infty)v(\xi) + Q(\xi)v(\xi),$$

Outline of proof: Essential spectrum of \mathcal{L}

Several planar polar coordinates:

$$\begin{aligned} [\mathcal{L}_2 v](\xi) = & A \left[\sum_{l=1}^k \left(\partial_{r_l}^2 + \frac{1}{r_l} \partial_{r_l} + \frac{1}{r_l^2} \partial_{\phi_l}^2 \right) + \sum_{l=2k+1}^d \partial_{x_l}^2 \right] v(\xi) \\ & - \sum_{l=1}^k \sigma_l \partial_{\phi_l} v(\xi) + Df(v_\infty) v(\xi) + Q(\xi) v(\xi), \end{aligned}$$

$$\xi = (r_1, \phi_1, \dots, r_k, \phi_k, x_{2k+1}, \dots, x_d), \quad Q(\xi) := Q(T_1(T_2(\xi)))$$

3. Simplified operator (far-field linearization): Neglecting $\mathcal{O}(\frac{1}{r})$ -terms yields

$$[\mathcal{L}_2^{\text{sim}} v](x) = A \left[\sum_{l=1}^k \partial_{r_l}^2 + \sum_{l=2k+1}^d \partial_{x_l}^2 \right] v(\xi) - \sum_{l=1}^k \sigma_l \partial_{\phi_l} v(\xi) + Df(v_\infty) v(\xi).$$

Outline of proof: Essential spectrum of \mathcal{L}

Simplified operator (far-field linearization):

$$[\mathcal{L}_2^{\text{sim}} v](\xi) = A \left[\sum_{l=1}^k \partial_{r_l}^2 + \sum_{l=2k+1}^d \partial_{x_l}^2 \right] v(\xi) - \sum_{l=1}^k \sigma_l \partial_{\phi_l} v(\xi) + Df(v_\infty) v(\xi)$$

4. Angular Fourier decomposition:

$$v(\xi) = \exp\left(i\omega \sum_{l=1}^k r_l\right) \exp\left(i \sum_{l=1}^k n_l \phi_l\right) \hat{v}, \quad n_l \in \mathbb{Z}, \omega \in \mathbb{R}, \hat{v} \in \mathbb{C}^N, |\hat{v}| = 1$$
$$\phi_l \in]-\pi, \pi], r_l > 0, l = 1, \dots, k,$$

yields

$$[(\lambda I - \mathcal{L}_2^{\text{sim}}) v](\xi) = \left(\lambda I_N + \omega^2 A + i \sum_{l=1}^k n_l \sigma_l I_N - Df(v_\infty) \right) v(\xi).$$

Outline of proof: Essential spectrum of \mathcal{L}

Angular Fourier decomposition:

$$[(\lambda I - \mathcal{L}_2^{\text{sim}}) v](\xi) = \left(\lambda I_N + \kappa^2 A + i \sum_{l=1}^k n_l \sigma_l I_N - Df(v_\infty) \right) v(\xi).$$

$$n_l \in \mathbb{Z}, \quad \kappa \in \mathbb{R}, \quad \pm i \sigma_l \text{ nonzero eigenvalues of } S \in \mathbb{R}^{d,d}$$

5. Finite-dimensional eigenvalue problem: $[(\lambda I - \mathcal{L}_2^{\text{sim}}) v](\xi) = 0$ for every ξ if $\lambda \in \mathbb{C}$ satisfies

$$(\omega^2 A - Df(v_\infty)) \hat{v} = - \left(\lambda + i \sum_{l=1}^k n_l \sigma_l \right) \hat{v}, \text{ for some } \omega \in \mathbb{R}.$$

Example

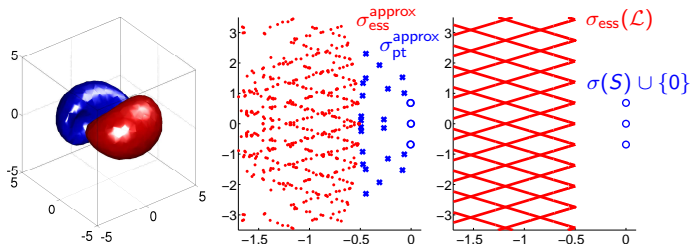
Consider the **quintic complex Ginzburg-Landau equation** (QCGL):

$$u_t = \alpha \Delta u + u \left(\mu + \beta |u|^2 + \gamma |u|^4 \right), \quad u = u(x, t) \in \mathbb{C}$$

with $u : \mathbb{R}^d \times [0, \infty[\rightarrow \mathbb{C}$, $d \in \{2, 3\}$. For the parameters

$$\alpha = \frac{1}{2} + \frac{1}{2}i, \beta = \frac{5}{2} + i, \gamma = -1 - \frac{1}{10}i, \mu = -\frac{1}{2}$$

this equation exhibits so called **spinning soliton** solutions.



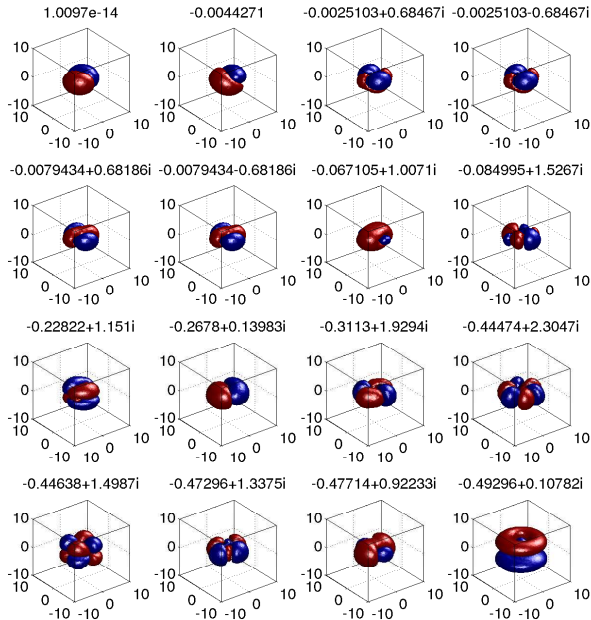


Figure: Eigenfunctions of QCGL for a spinning soliton with $d = 3$

Work in progress

- exponential decay in **space of continuous functions**
- rotating waves in **bounded domains**
- **approximation theorem** for rotating waves
- asymptotic boundary conditions
- **numerical computations** (interaction of multisolitons)

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